

# On the Collapsing Sandpile Problem

Serge DUMONT, Nouredine IGBIDA\*

March 2008

## Abstract

In this note, we are interested in the numerical analysis of the collapse of an unstable sandpile. By using the collapsing model introduced by Evans, Feldman and Garipey in [6], we give a description of the phenomena in terms of a composition of projections onto interlocked convex sets around the set of stable sandpiles.

## 1. Introduction and main results

We are interested in the collapse of an unstable sandpile; i.e. taking some sandpile with an unstable initial configuration, what is the final resting state of the sandpile after various avalanches. To describe the problem we need to use some function  $u : \mathbb{R}^D \rightarrow \mathbb{R}$ , the height function of the sandpile, where  $D \geq 1$  (in practice  $D = 2$ ). The stability constraint of the sandpile reads

$$(1) \quad |\nabla u| \leq 1,$$

and has the physical meaning that the sand cannot remain in equilibrium if the slope anywhere exceeds angle  $\pi/4$ . Now, if  $g$  represents the height of a starting unstable profile (assume for instance  $g$  is the profile of a wet sandpile) i.e.

$$L := \|\nabla g\|_\infty > 1,$$

then constraint (1) forces the height function rearranges itself to achieve the profile in a stable configuration. Our main interest lies into the description of mapping  $\mathbb{Q} : g \rightarrow \mathbb{Q}(g)$ , where  $\mathbb{Q}(g)$  is the final stable profile of the sandpile associated with the initial unstable one  $g$ .

---

\*LAMFA, UPJV CNRS UMR 6140, Université de Picardie Jules Verne, 33 rue Saint Leu, 80038 Amiens, France. Email : [serge.dumont,nouredine.igbida]@u-picardie.fr

By using the Monge Kantorovich mass transfer theory, Evans, Feldman and Gariépy [6] introduced a simplistic model for the collapse of an unstable sandpile. It is given by the limit of the flow governed by the  $p$ -Laplacian, as  $p \rightarrow \infty$  :

$$(2) \quad \begin{cases} -u_t = \Delta_p u & \text{in } Q := (0, \infty) \times \Omega \\ u(0) = g & \text{in } \Omega. \end{cases}$$

Letting  $p \rightarrow \infty$ , one expects that the limit of solution  $u_p$  satisfies evolution equation (multivalued)

$$(3) \quad u_t \in \partial \mathbb{I}_K(u) \quad \text{in } Q,$$

where  $\mathbb{I}_K$  is defined by

$$\mathbb{I}_K(u) = \begin{cases} 0 & \text{if } u \in K \\ +\infty & \text{otherwise,} \end{cases}$$

$$K = \left\{ u \in W^{1,\infty}(\Omega) ; |\nabla u| \leq 1 \text{ a.e. } \Omega \right\}.$$

and the subdifferential operator  $\partial \mathbb{I}_K$  is defined in  $L^2(\Omega)$ , by

$$h \in \partial \mathbb{I}_K(u) \Leftrightarrow h \in L^2(\Omega), u \in K \text{ and } \int_{\Omega} h(z-u) \leq 0 \quad \forall z \in L^2(\Omega).$$

It is clear that compatible initial data for (3) leaves in  $K$ , so that the limit of the solution of (2), when  $g \notin K$ , is singular: letting  $p \rightarrow \infty$ , turns out to grow up a boundary layer connecting  $g$  to the limit. In [6] (see also [3]), it is proved that the limiting function is  $v(1)$  (independent of  $t$ ), where  $v$  is the unique solution of

$$(4) \quad \begin{cases} v(t)/t - v_t(t) \in \partial \mathbb{I}_{\infty}(v(t)) & \text{a.e. } t \in (\delta, 1] \\ v(\delta) = \delta g \end{cases}$$

and  $\delta = 1/\|\nabla g\|_{\infty}$ , in the sense that  $v \in W^{1,2}(\delta, 1; L^2(\Omega))$ ,  $v(\delta) = \delta g$  and, for any  $t \in (\delta, 1]$ ,  $v(t)/t - v_t(t) \in \partial \mathbb{I}_{\infty}(v(t))$ .

In terms of Monge Kantorovich mass transfer theory the equation (4) means that  $v$  is a potential corresponding to optimal moving the mass  $\mu^+ = v(\cdot, t)/t dx$  to  $\mu^- = v_t(\cdot, t) dx$ .

As a consequence of [6] (see also [3]), we have the following characterization of mapping  $\mathbb{Q}$  in terms of evolution equation (4):

**Theorem 1** ([6] and [3]) *For any  $g \in L^2(\Omega)$ ,  $\mathbb{Q}(g) = v(1)$ , where  $v$  is the unique solution of (4).*

In [6], it is also proved that operator  $\mathbb{Q}$  is not the projection onto  $K$ , the set of stable sandpiles. In order to give the numerical analysis and simulation of the collapsing of a sandpile, we give, in this paper, a more precise description of  $\mathbb{Q}$ , in terms of a composition of projections onto interlocked convex sets of  $L^2(\Omega)$ , around convex  $K$ .

To set our main results, let us give some notations. For a given convex  $C \subseteq L^2(\Omega)$ , we denote by  $\mathbb{P}_C$ , the projection with respect to the  $L^2$  norm on  $C$ , defined by:

$$z = \mathbb{P}_C(u) \Leftrightarrow z \in C, \int_{\Omega} (u - z)(v - z) \leq 0 \quad \text{for any } v \in C.$$

For a given  $[a, b]$  compact interval of  $\mathbb{R}$ , we say that  $(d_i)_{i=0}^n$  is an  $\varepsilon$ -discretization of  $(a, b)$ , provided  $\varepsilon > 0$ ,  $d_0 = a < d_1 < d_2 < \dots < d_n = b$  and  $d_i - d_{i-1} \leq \varepsilon$ , for any  $i = 1, \dots, n$ . Notice that, letting  $\varepsilon \rightarrow 0$  in an  $\varepsilon$ -discretization of a fixed interval  $(a, b)$ , is equivalent to let  $n \rightarrow \infty$ . For any  $d > 0$ , we denote by  $K(d)$ , the convex set given by

$$K(d) = \left\{ z \in W^{1,\infty}(\Omega) ; |\nabla z| \leq d \right\}.$$

**Theorem 2** *Let  $g \in W^{1,\infty}(\Omega)$ ,  $a := \|\nabla g\|_{\infty}$  and  $(d_i)_{i=0}^n$  an  $\varepsilon$ -discretization of  $(1, a)$ . Then*

$$(5) \quad \mathcal{Q}(g) = \lim_{n \rightarrow \infty} \mathbb{P}_{K(1)} \mathbb{P}_{K(d_1)} \dots \mathbb{P}_{K(d_{n-2})} \mathbb{P}_{K(d_{n-1})} g.$$

**Proof :** For  $i = 1, \dots, n$ , let us denote by

$$u_i = \mathbb{P}_{K(d_{n-i})} u_{i-1} \quad \text{and } u_0 = g,$$

i.e.

$$u_i + \partial \mathbb{I}_{K(d_{n-i})}(u_i) \ni u_{i-1} \quad \text{for } i = 1, \dots, n,$$

and

$$\mathbb{P}_{K(1)} \mathbb{P}_{K(d_1)} \dots \mathbb{P}_{K(d_{n-2})} \mathbb{P}_{K(d_{n-1})} g = u_n.$$

Setting, for  $i = 1, \dots, n$ ,

$$z_i = u_i / d_{n-i},$$

it is not difficult to see that

$$(6) \quad z_i + \partial \mathbb{I}_{K(1)}(z_i) \ni \frac{d_{n-i+1}}{d_{n-i}} z_{i-1}.$$

Now, setting

$$t_i = 1/d_{n-i} \quad \text{for } i = 0, \dots, n,$$

it is clear that  $(t_i)_{i=0}^n$  is an  $\varepsilon$ -discretization of  $(\delta, 1)$ . So that, the  $\varepsilon$ -approximate solution of

$$\begin{cases} v_i(t) + \partial \mathbb{I}_{\infty}(v(t)) \ni f(t) & \text{a.e. } t \in (\delta, 1] \\ v(\delta) = \delta g, \end{cases}$$

with  $f(t) = v(t)/t$ , associated with  $(t_i)_{i=0}^n$  converges to  $v$  in  $\mathcal{C}([\delta, 1]; L^2(\Omega))$ . More precisely, taking the Euler implicit discretization in time

$$(7) \quad v_i + \partial \mathbb{I}_{K(1)}(v_i) \ni v_{i-1} + \frac{t_i - t_{i-1}}{t_{i-1}} v(t_{i-1}), \quad i = 1, 2, \dots, n,$$

where, we used the discretization of  $f$  given by  $f_i = \frac{v(t_{i-1})}{t_{i-1}}$ , for  $i = 1, \dots, n$ , and defining  $\varepsilon$ -approximate solution  $v_\varepsilon$  by  $v_\varepsilon(t) = v_i$  for  $t \in [t_i, t_{i+1}]$  for  $i = 0, 1, \dots, n-1$ , we have

$$(8) \quad v_\varepsilon \rightarrow v \quad \text{in} \quad \mathcal{C}([\delta, 1]; L^2(\Omega)), \text{ as } \varepsilon \rightarrow 0,$$

and

$$(9) \quad v_n \rightarrow v(1) \quad \text{in} \quad L^2(\Omega), \text{ as } n \rightarrow \infty.$$

Moreover, it is not difficult to see that (7) is equivalent to

$$v_i + \partial \mathbb{I}_{K(1)}(v_i) \ni \frac{d_{n-i+1}}{d_{n-i}} v_{i-1} + \frac{t_i - t_{i-1}}{t_{i-1}} (v(t_{i-1}) - v_{i-1}),$$

so that, by using (6) and the  $L^2$ -contraction property of  $(I + \partial \mathbb{I}_{K(1)})^{-1}$ , for  $i = 1, \dots, n$ , we get

$$(10) \quad \begin{aligned} \|v_i - z_i\|_2 &\leq \frac{d_{n-i+1}}{d_{n-i}} \|v_{i-1} - z_{i-1}\|_2 + \frac{t_i - t_{i-1}}{t_{i-1}} \|v(t_{i-1}) - v_{i-1}\|_2 \\ &\leq \frac{t_i}{t_{i-1}} \|v_{i-1} - z_{i-1}\|_2 + \frac{t_i - t_{i-1}}{t_{i-1}} \|v(t_{i-1}) - v_{i-1}\|_2. \end{aligned}$$

Since,  $v_0 = z_0 = g$ , then iterating (10) for  $i = n, \dots, 1$ , we obtain

$$\begin{aligned} \|v_n - z_n\|_2 &\leq \sum_{i=2}^n \frac{t_n}{t_{n-i+1}} \frac{t_{n-i+1} - t_{n-i}}{t_{n-i}} \|v(t_{n-i}) - v_{n-i}\|_2 + \frac{t_n - t_{n-1}}{t_{n-1}} \|v(t_{n-1}) - v_{n-1}\|_2 \\ &\leq \frac{1}{\delta^2} \sum_{i=2}^{n-1} (t_{n-i+1} - t_{n-i}) \|v(t_{n-i}) - v_{n-i}\|_2 + \frac{1}{\delta} (t_n - t_{n-1}) \|v(t_{n-1}) - v_{n-1}\|_2 \\ &\leq \frac{1}{\delta^2} \sum_{i=0}^{n-1} (t_{n-i+1} - t_{n-i}) \|v(t_{n-i}) - v_{n-i}\|_2 \\ &\leq \frac{1}{\delta^2} \sum_{i=1}^{n-1} (t_{i+1} - t_i) \|v(t_i) - v_i\|_2, \end{aligned}$$

where we used the fact that, for any  $k = 0, \dots, n$ ,  $\delta \leq t_k \leq 1$  and  $1/t_k \leq 1/\delta \leq 1/\delta^2$ . Considering  $\bar{v}_\varepsilon$  given by  $\bar{v}_\varepsilon(t) = v(t_i)$  for  $t \in [t_i, t_{i+1}[$  and  $i = 0, 2, \dots, n-1$ , we obtain

$$\begin{aligned} \|v_n - z_n\|_2 &\leq \frac{1}{\delta^2} \sum_{i=0}^{n-1} \int_{t_i}^{t_{i+1}} \|\bar{v}_\varepsilon(t) - v_\varepsilon(t)\|_2 dt \\ &\leq \frac{1}{\delta^2} \int_\delta^1 \|\bar{v}_\varepsilon(t) - v_\varepsilon(t)\|_2 dt \end{aligned}$$

Since, as  $\varepsilon \rightarrow 0$ ,  $\bar{v}_\varepsilon \rightarrow v$  and  $v_\varepsilon \rightarrow v$  in  $\mathcal{C}([\delta, 1]; L^2(\Omega))$ , then

$$\lim_{n \rightarrow \infty} \|v_n - z_n\|_2 = 0.$$

At last, since  $z_n = u_n/d_0 = u_n$ , then by using (9), we deduce that, as  $n \rightarrow \infty$ ,

$$u_n \rightarrow v(1) \quad \text{in} \quad L^2(\Omega),$$

and (5) follows. ■

## 2. Numerical approximations

### 2.1 The numerical problem

Let  $g \in W^{1,\infty}(\Omega)$  be given. Our aim in this section is to give a numerical approximation of  $\mathbb{Q}(g)$ . By using Theorem 2, it is clear that it is enough to find an numerical method to compute the projection of a given function, that we call  $g$  again, on a convex  $K_{(d)}$  with arbitrary  $d \geq 1$ .

In [5], we gave a numerical method based on duality arguments for numerical approximation of the projection on the convex  $K_{(1)}$ . Even if the context here is different, we are going to use the same method for our situation. Remember that

$$(11) \quad u = \mathbb{P}_{K_{(d)}}(g)$$

is equivalent to

$$J(u) = \inf\{J(v), v \in W^{1,\infty}(\Omega)\}$$

where  $J(v) = \frac{1}{2}\|v - g\|_{L^2}^2 + \mathbb{I}_{K_{(d)}}(v)$ . In, [5], we have proved that:

$$(12) \quad \min_{v \in W^{1,\infty}(\Omega)} J(v) = \sup_{q \in H_{div}(\Omega)} -G(q)$$

where

$$G(q) = \frac{1}{2} \int_{\Omega} |div q|^2 + \int_{\Omega} g div q + \int_{\Omega} |q|$$

and

$$H_{div}(\Omega) = \{q \in (L^2(\Omega))^2, div(q) \in L^2(\Omega)\}.$$

Remember also, that in general  $\sup_{q \in H_{div}(\Omega)} -G(q)$  is not reached in  $H_{div}(\Omega)$ . Indeed,

there exists  $\sigma$  a  $\mathbb{R}^N$ -vector valued Radon measure such that  $div(\sigma) \in L^2(\Omega)$ ,  $u = \mathbb{P}_{K_{(d)}}g = g - div\sigma$  in  $\mathcal{D}'(\Omega)$  and  $\sigma$  is related to the tangential derivative of  $u$  with respect to  $|\sigma|$ , the total variation measure associated to  $\sigma$  (cf. [4]). For the exact description of  $\sigma$  one can see the paper [11] for equivalent formulations in divergence form of equation  $v \in \partial \mathbb{I}_{K_{(d)}}f$ .

So, an numerical approximation of  $\mathbb{P}_{K_{(d)}}g$  will follow by approximating  $\sup_{q \in H_{div}(\Omega)} -G(q)$  and computing  $g - div(q)$ . Coming back to the numerical approximation of  $\mathbb{Q}(g)$ , we take  $(d_i)_{i=0}^n$  an  $\varepsilon$ -discretization of interval  $(1, \|g\|_{\infty})$ , with a given small  $\varepsilon > 0$  ; i.e. a large  $n$ , and we use the algorithm of approximation of  $\mathbb{P}_{K_{(d)}}g$ , to approximate composition

$$\mathbb{P}_{K_{(1)}}\mathbb{P}_{K_{(d_1)}}\dots\mathbb{P}_{K_{(d_{n-2})}}\mathbb{P}_{K_{(d_{n-1})}}g.$$

## 2.2 Space discretization

The solution of dual problem (12) is computed using Raviart Thomas finite element of the lowest order [12]. Denoting by  $h$  the average length of the elements,  $V_h$  the space of finite elements and  $\sigma_h$  the solution of problem

$$(13) \quad G(\sigma_h) = \inf_{q_h \in V_h} G(q_h),$$

the convergence of  $\sigma_h$  to  $\sigma$  is guaranteed remarking that we have

$$\|\sigma_h - \sigma\|_{L^2} \leq \|r_h(\sigma_p) - \sigma\|_{L^2}$$

where  $r_h$  is the projection onto  $V_h$ .

Finally, the minimization of  $G$  on  $V_h$  is implemented using a relaxation procedure.

## 3. Numerical results

In this section we present some numerical results on collapsing sandpiles.

### 3.1 The one dimensional case

First, we study example given in [6], where the stable sandpile is not the projection of the initial data. More precisely, the unstable initial data is composed in the following manner: two sandpiles with sides of slope equal to 2, with a width equal to 0.2, and centered around  $x_0 = \pm 0.22$ ; and one sandpile with sides of slope equal to 250, with a width equal to 0.008 and centered around  $x = 0$ . The spacial discretization of the problem is carried out using  $N = 2500$  points.

The maximal slope of this sandpile is  $|\nabla g|_\infty = \bar{\theta} = 250$ .

Figure 1 presents solutions obtained using different uniform discretizations of interval  $[1, \bar{\theta}]$ . In others words, decreasing maximal slope  $\theta_i$  for sequence of projections is chosen such that  $\theta_i = 1 + i \frac{\bar{\theta}-1}{M}$  for  $i = 1, \dots, M$ . In the numerical experiments, the number  $M$  of projections is ranging from 1 to 5000.

One can observe that if we do a direct projection on  $K(1)$  ( $M = 1$ ), then the numerical final stable sandpile is composed of three adjacent sandpiles. When number  $M$  of subdivisions increases, the central small triangle disappears.

These numerical results corroborate the theoretical results given in [6], where it is shown that in this case, the stable sandpile consists of two adjacent triangles.

### 3.2 The two dimensional case

This paragraph is devoted to the two dimensional computations.

First, figure 3 shows the convergence rate of the approximation of one projection, when the average step of discretization tends to zero. In this computation, domain  $\Omega$  is unit square  $(-1, 1)^2$ , discretized with squares  $[-1 + ih, -1 + (i + 1)h] \times [-1 + jh, -1 + (j + 1)h]$ , with  $h = \frac{2}{N}$  and  $0 \leq i, j \leq N$  for a given integer  $N$ .

Initial data  $g$  is a cone centered on the origin with radius of the basis equal to  $R = \frac{1}{4}$ , height  $H = 1$  (see figure 2, up), and exact solution  $u$  is a cone centered on the

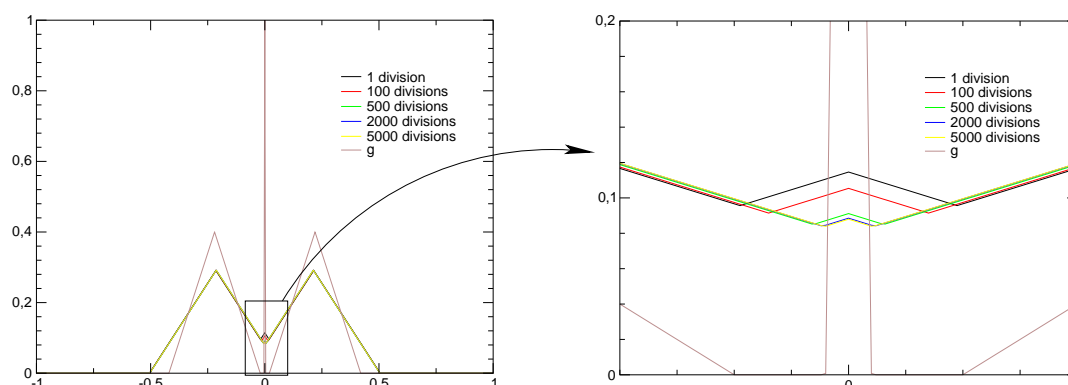


Figure 1: The counterexample where the final stable sandpile is not the direct projection of the initial unstable sandpile (details on the right)

origin with radius of the basis  $\tilde{R}$  and height  $\tilde{H}$  equal to  $\tilde{R} = \tilde{H} = 2^{-\frac{4}{3}} \simeq 0.397$  (see figure 2, down).

Results plotted in figure 3 shows the  $L^2$ -norm of the error  $\|u_N - u\|_{L^2(\Omega)}$  where  $u_N$  is the computed solution, versus number  $N = \frac{2}{h}$  of squares used for the discretization in each direction. This figure shows that the convergence rate of the error verifies  $\|u_N - u\|_{L^2(\Omega)} = O(N^{-\alpha})$  with  $\alpha \simeq 1.92$ , when  $N$  is ranging from 50 to 150.

We present here a result with a more complex initial unstable sandpile, plotted in figure 4. The maximal slope of this sandpile is equal to  $|\nabla g|_\infty = 4.1$ , and the maximal slope of the final stable sandpile is equal to  $|\nabla u|_\infty = 0.4$ .

**Remarks** *In the theoretical part of this article, for seek of simplicity, angle  $\alpha$  of the maximal slope of stability has been taken equal to  $\alpha = \frac{\pi}{4}$  (see formula (1) for example), but the results presented here are true with any angle  $\alpha \in (0, \frac{\pi}{2})$ . Here, we have chosen  $\alpha \simeq \frac{\pi}{8}$ .*

Two different boundary conditions are considered to compute the projection:

- First, we take the Dirichlet boundary condition:  $u(x) = 0$  on  $\partial\Omega$ . This simulates the experiment where the sandpile collapses on a table. The result is plotted in figure 5.
- Then, we consider the Neumann boundary condition:  $m\partial_n u = 0$  on  $\partial\Omega$ . This simulates the case where domain  $\Omega$  is the bottom of a box. Here, the mass has to be conserved. The final stable sandpile is plotted in figure 6. The initial and the final sandpile has a volume equal to 0.62178, which shows that the volume is numerically conserved.

These results are obtained using a spacial discretization with  $N = 100$  and a number of emboities projections equal to  $M = 40$ . In this case, the result does not depend on  $M$ .

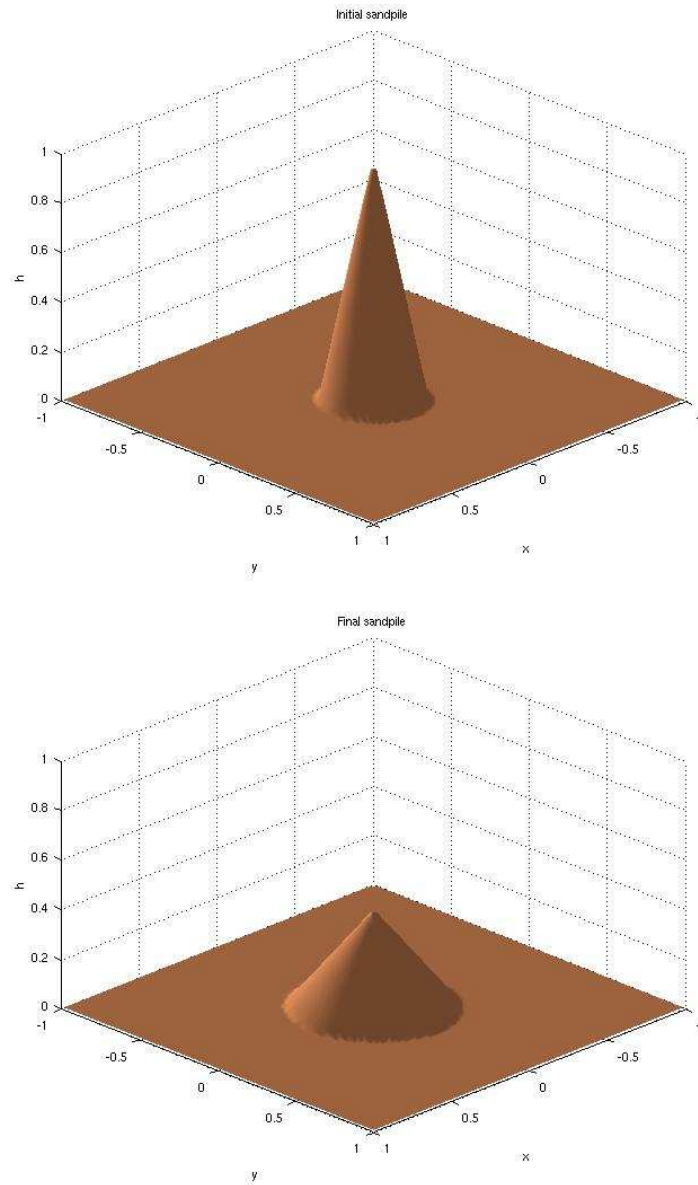


Figure 2: Initial unstable (up) and final stable (down) sandpile



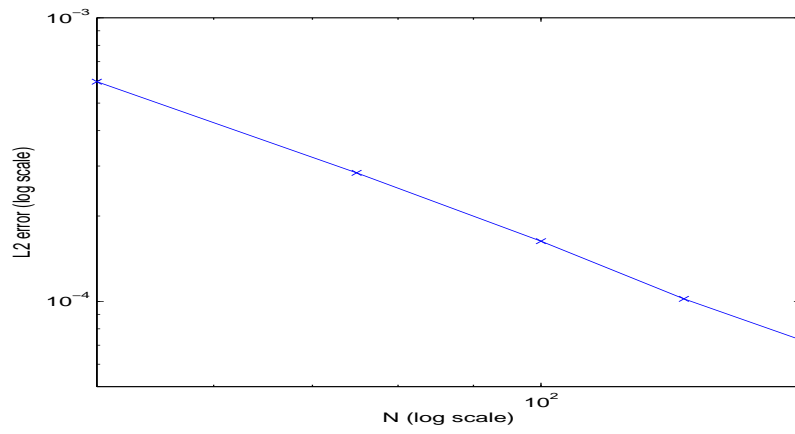


Figure 3: Accuracy of the method ( $L^2$ -norm of the difference between the exact and the computed solution)

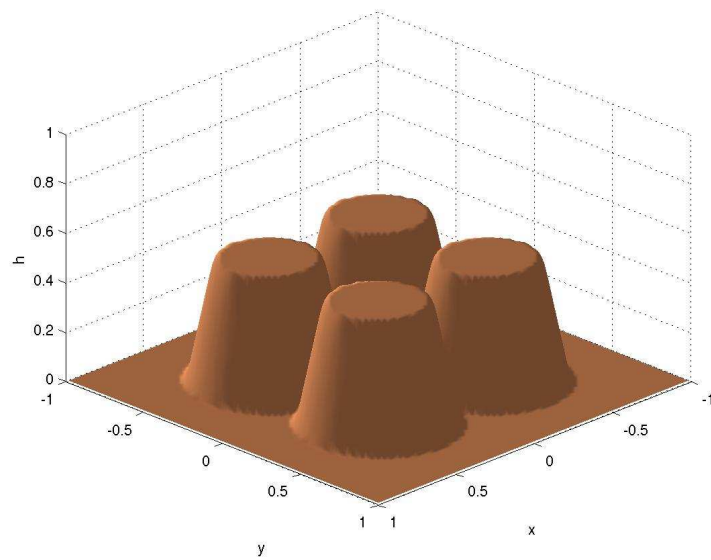


Figure 4: Initial unstable sandpile

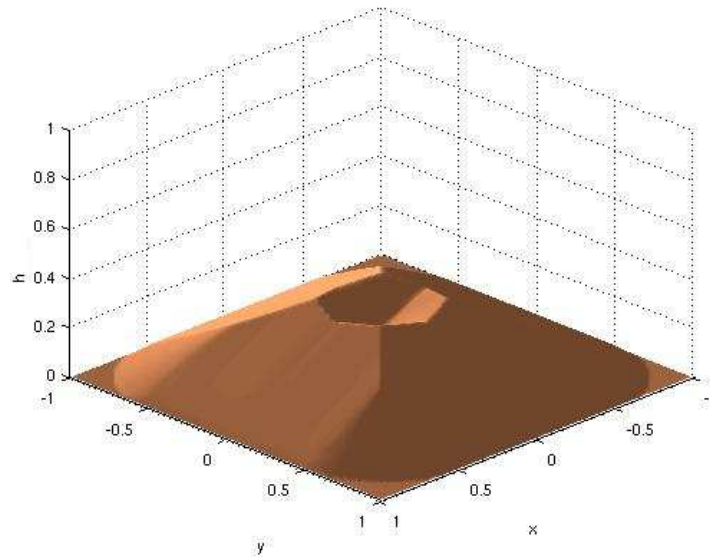


Figure 5: Collapsing sandpile with Dirichlet boundary condition

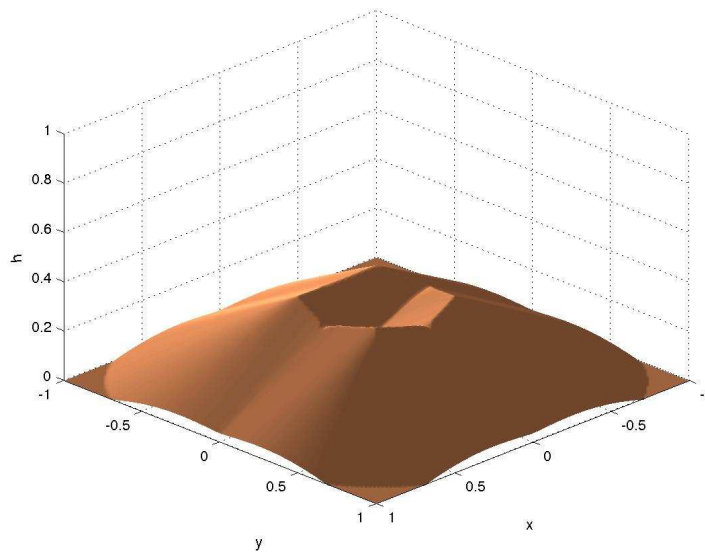


Figure 6: Collapsing sandpile with Neumann boundary condition

## References

- [1] L. De Pascale, L. C. Evans and A. Pratelli, Integral estimates for transport densities, *Bull. London Math. Soc.*, **36**, 3, 383-395, 2004.
- [2] J. W. Barrett and L. Prigozhin, Dual formulation in Critical State Problems, *Interfaces and Free Boundaries*, **8**, 349-370, 2006.
- [3] Ph. Bénilan, L. C. Evans and R. F. Gariepy, On some singular limits of homogeneous semigroups, *J. Evol. Equ.*, **3** (2003), no. 2, 203–214.
- [4] G. Bouchitté, G. Buttazzo and P. Seppecher, Energies with respect to a Measure and Applications to Low Dimensional Structures, *Calc. Var. Partial Differential Equations* 5 (1997), 37-54.
- [5] S. Dumont and N. Igbida, Back on a Dual Formulation for the Growing Sandpile Problem, *Preprint*, 2007.
- [6] L. C. Evans, M. Feldman and R. F. Gariepy, Fast/Slow diffusion and collapsing sandpiles, *J. Differential Equations*, **137**:166–209, 1997.
- [7] L. C. Evans and W. Gangbo, Differential equations methods for the Monge-Kantorovich mass transfer problem, *Mem. Am. Math. Soc.*, **137**, no. 653, 1999.
- [8] I. Ekeland and R. Témam, Convex analysis and variational problems, *Classics in Applied Mathematics, Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1999.*
- [9] L. C. Evans, J.-L. Lions et R. Trémollières, Partial differential equations and Monge-Kantorovich mass transfer, *Current Developments in Mathematics*, Int. Press, Boston, Ma, (1997) 65-126.
- [10] R. Glowinski, J.-L. Lions and R. Trémollières, Analyse numérique des inéquations variationnelles, *Méthodes Mathématiques de l'Informatique, Dunod, Paris, 1976.*
- [11] N. Igbida, On Monge-Kantorovich Equation, *Preprint*, 2008.
- [12] J.E. Roberts and J.-M. Thomas, Mixed and hybrid methods, P.G. Ciarlet, J.L. Lions (Eds.), *Handbook of Numerical Analysis*, vol. II, Finite Element Methods (Part 1), North-Holland, Amsterdam, 1991.